

PAPER • OPEN ACCESS

On the two-component generalization of the (2+1)-dimensional Davey-Stewartson I equation

To cite this article: Nurzhan Serikbayev *et al* 2019 *J. Phys.: Conf. Ser.* **1391** 012160

View the [article online](#) for updates and enhancements.

You may also like

- [Soliton resonance and web structure in the Davey-Stewartson system](#)
Gino Biondini, Dmitri Kireyev and Ken-ichi Maruno
- [Global existence and asymptotic behaviour in time of small solutions to the elliptic - hyperbolic Davey - Stewartson system](#)
Nakao Hayashi and Hitoshi Hirata
- [Numerical scattering for the defocusing Davey-Stewartson II equation for initial data with compact support](#)
Christian Klein and Nikola Stoilov



The Electrochemical Society
Advancing solid state & electrochemical science & technology

242nd ECS Meeting

Oct 9 – 13, 2022 • Atlanta, GA, US

Early hotel & registration pricing
ends September 12

Presenting more than 2,400
technical abstracts in 50 symposia

The meeting for industry & researchers in

BATTERIES
ENERGY TECHNOLOGY
SENSORS AND MORE!



Register now!



**ECS Plenary Lecture featuring
M. Stanley Whittingham,**
Binghamton University
Nobel Laureate –
2019 Nobel Prize in Chemistry



On the two-component generalization of the (2+1)-dimensional Davey-Stewartson I equation

Nurzhan Serikbayev¹, Gulgassyl Nugmanova², Ratbay Myrzakulov³

¹³ L.N.Gumilyov Eurasian National University, General and Theoretical Physics, Kazakhstan

² L.N.Gumilyov Eurasian National University, Mathematical and Computer Modeling, Kazakhstan

E-mail: ¹ns.serikbayev@gmail.com, ²nugmanovagn@gmail.com, ³rmyrzakulov@gmail.com

Abstract. The geometric-gauge equivalent of the famous Ishimori spin equation is the (2+1)-dimensional Davey-Stewartson equation, which in turn is one of the (2+1)-dimensional generalizations of the nonlinear Schrodinger equation. Multicomponent generalization of nonlinear integrable equations attract considerable interest from both physical and mathematical points of view. In this paper, the two-component integrable generalization of the (2+1)-dimensional Davey-Stewartson I equation is obtained based on its one-component representation, and the corresponding Lax representation is also obtained.

1. Introduction

It is known that integrable nonlinear Schrodinger type equations (NSE) are key models in the theory of integrable equations. Recently, their multicomponent generalizations have been actively studied. In the work [1] it is shown that the Manakov two-step system is integrable. The geometrical connection with the last system and the two-layer spin model was established in [2] - [4]. The geometric gauge equivalent of the Ishimori spin equation [5] is a (2+1)-dimensional Davy-Stewartson (DS) equation, which is one of the (2+1) -dimensional generalizations of the NSE [6].

Consider the (2+1)-dimensional DS equation

$$iq_t + \frac{1}{2}(\sigma^2 q_{xx} + q_{yy}) = (v - qr)q, \quad (1)$$

$$-ir_t + \frac{1}{2}(\sigma^2 r_{xx} + r_{yy}) = (v - qr)r, \quad (2)$$

$$v_{xx} - \sigma^2 v_{yy} = 2(qr)_{xx}, \quad (3)$$

where $r = \pm q^*$, q^* is the complex conjugate of q . For $\sigma^2 = 1$, the system of equations 1-2 is called the Davy-Stewartson Type I equation (DSI), and for $\sigma^2 = -1$ – the Davy-Stewartson Type II equation (DSII) [7]. We focus on the (2+1)-dimensional DSI equation, and in the next section we give some well-known data on the single-component (2+1)-dimensional DSI equation. Our goal is to derive a two-component generalization for the last equation.

The result of the study is formed in the form of approval and is proved in the second paragraph.



2. A single-component (2+1)-dimensional DSI equation

As it is known, the standard Lax representation of the DSI equation is

$$F_y = \sigma_3 F_x + QF, \quad (4)$$

$$F_t = A_2 F_{xx} + A_1 F_x + A_0 F, \quad (5)$$

where

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \begin{pmatrix} 0 & q \\ -r & 0 \end{pmatrix},$$

$$A_0 = i \begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix}, \quad A_1 = 2i \begin{pmatrix} 0 & q \\ -q & 0 \end{pmatrix}, \quad A_2 = 2i\sigma_3.$$

The elements of the matrix A_0 satisfy the following conditions:

$$c_{12} = \frac{1}{2}(\partial_x + \partial_y)q,$$

$$c_{21} = -\frac{1}{2}(\partial_x - \partial_y)r,$$

$$(\partial_x - \partial_y)c_{11} = -\frac{1}{2}(\partial_x + \partial_y)(qr),$$

$$(\partial_x + \partial_y)c_{22} = \frac{1}{2}(\partial_x - \partial_y)(qr),$$

here $\partial_x \equiv \partial/\partial x$, $\partial_y \equiv \partial/\partial y$ and the field v in the system of DS equations (1)-(2) is defined as

$$v = -i(c_{22} - c_{11}) + qr.$$

The compatibility condition for the equation (4) and (5) $F_{yt} = F_{ty}$ implies the following series of equations:

$$[\sigma_3, A_2] = 0, \quad (6)$$

$$\sigma_3 A_{2x} - A_{2y} + [\sigma_3, A_1] + [Q, A_2] = 0, \quad (7)$$

$$\sigma_3 A_{1x} - A_{1y} + [\sigma_3, A_0] + [Q, A_1] - 2A_2 Q_x = 0, \quad (8)$$

$$\sigma_3 A_{0x} - A_{0y} + Q_t + [Q, A_0] - A_2 Q_{xx} - A_1 Q_x = 0. \quad (9)$$

From the system of equations (6)-(9) it is not difficult to obtain a one-component (2+1)-dimensional DSI equation [8].

3. A two-component (2+1)-dimensional DSI equation

Theorem. *If the matrices Σ , Q and A belong to the group $SU(3)$, then a two-component (2+1)-dimensional DSI equation has the following form:*

$$iq_{1t} + q_{1xx} + q_{1yy} - v_1 q_1 - w_1 q_2 = 0, \quad (10)$$

$$iq_{2t} + q_{2xx} + q_{2yy} - w_2 q_1 - v_2 q_2 = 0, \quad (11)$$

$$-ir_{1t} + r_{1xx} + r_{1yy} - v_1 r_1 - w_1 r_2 = 0, \quad (12)$$

$$-ir_{2t} + r_{2xx} + r_{2yy} - w_2 r_1 - v_2 r_2 = 0, \quad (13)$$

$$v_{1xx} - v_{1yy} = (2r_1 q_1 + r_2 q_2)_{xx} + 2(r_2 q_2)_{xy} + (2r_1 q_1 + r_2 q_2)_{yy}, \quad (14)$$

$$v_{2xx} - v_{2yy} = (r_1 q_1 + 2r_2 q_2)_{xx} + 2(r_1 q_1)_{xy} + (r_1 q_1 + 2r_2 q_2)_{yy}, \quad (15)$$

$$w_{1xx} - w_{1yy} = (q_1 r_2)_{xx} - 2(q_1 r_2)_{xy} + (q_1 r_2)_{yy}, \quad (16)$$

$$w_{2xx} - w_{2yy} = (q_2 r_1)_{xx} - 2(q_2 r_1)_{xy} + (q_2 r_1)_{yy}, \quad (17)$$

where q and r are complex-valued functions, and v_j and w_j are real functions.

Proof. For proof, we require that the column matrix F satisfies the following Lax representation

$$F_y = \Sigma F_x + PF, \quad (18)$$

$$F_t = B_2 F_{xx} + B_1 F_x + B_0 F, \quad (19)$$

where

$$\Sigma = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix},$$

the remaining matrices belong to the group $su(3)$:

$$P = \begin{pmatrix} 0 & q_1 & q_2 \\ -r_1 & 0 & 0 \\ -r_2 & 0 & 0 \end{pmatrix},$$

$$B_0 = \begin{pmatrix} c_{11} & c_{12} & c_{13} \\ c_{21} & c_{22} & c_{23} \\ c_{31} & c_{32} & c_{33} \end{pmatrix}, \quad B_1 = \begin{pmatrix} b_{11} & b_{12} & b_{13} \\ b_{21} & b_{22} & b_{23} \\ b_{31} & b_{32} & b_{33} \end{pmatrix}, \quad B_2 = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix}.$$

Then, from the compatibility condition $F_{yt} = F_{ty}$ of the system (18) and (19), we obtain

$$[\Sigma, B_2] = 0, \quad (20)$$

$$\Sigma B_{2x} - B_{2y} + [\Sigma, B_1] + [P, B_2] = 0, \quad (21)$$

$$\Sigma B_{1x} - B_{1y} + [\Sigma, B_0] + [P, B_1] - 2B_2 P_x = 0, \quad (22)$$

$$\Sigma B_{0x} - B_{0y} + P_t + [P, B_0] - B_2 P_{xx} - B_1 P_x = 0. \quad (23)$$

Now we define elements of the matrices B_0 , B_1 and B_2 . From the equation 20 it is determined that

$$a_{12} = a_{21} = a_{13} = a_{31} = 0.$$

Therefore, we have

$$B_2 = \begin{pmatrix} a_{11} & 0 & 0 \\ 0 & a_{22} & a_{23} \\ 0 & a_{32} & a_{33} \end{pmatrix}.$$

Similarly, from the equations (21) and (22) we get a number of restrictions for the elements of the matrices B_0 , B_1 and B_2 .

Namely, for the elements a_{ij} ($i, j = 1, 2, 3$) of the matrix B_2 :

$$a_{11x} + a_{11y} = 0,$$

$$a_{22x} + a_{22y} = 0,$$

$$a_{33x} + a_{33y} = 0,$$

$$a_{23x} + a_{23y} = 0,$$

$$a_{32x} + a_{32y} = 0,$$

for the elements b_{ij} ($i, j = 1, 2, 3$) of the matrix B_1 :

$$b_{12} = -\frac{1}{2}(a_{2211}q_1 + a_{32}q_2),$$

$$\begin{aligned}
b_{21} &= \frac{1}{2}(a_{2211}r_1 + a_{23}r_2), \\
b_{13} &= -\frac{1}{2}(a_{3311}q_2 + a_{23}q_1), \\
b_{31} &= \frac{1}{2}(a_{3311}r_2 + a_{32}r_1), \\
b_{11x} - b_{11y} &= 0, \\
b_{22x} + b_{22y} &= -\frac{1}{2}(a_{23}q_1r_2 - a_{32}q_2r_1), \\
b_{33x} + b_{33y} &= \frac{1}{2}(a_{23}q_1r_2 - a_{32}q_2r_1), \\
b_{23x} + b_{23y} &= \frac{1}{2}(a_{3322}q_2r_1 + a_{23}(q_1r_1 - q_2r_2)), \\
b_{32x} + b_{32y} &= -\frac{1}{2}(a_{3322}q_1r_2 + a_{32}(q_1r_1 - q_2r_2)),
\end{aligned}$$

where $a_{iijj} = a_{ii} - a_{jj}$ ($i > j$). And for the elements c_{ij} ($i, j = 1, 2, 3$) of the matrix B_0 , we have

$$\begin{aligned}
c_{12} &= \frac{1}{4}[(3a_{11} + a_{22})q_{1x} - a_{2211}q_{1y} + a_{32}q_{2x} - a_{32}q_{2y} + (a_{22x} - a_{22y} - 2b_{2211})q_1 + (a_{32x} - a_{32y} - 2b_{32})q_2], \\
c_{21} &= \frac{1}{4}[(3a_{11} + a_{22})r_{1x} - a_{2211}r_{1y} + 3a_{32}r_{2x} - a_{23}r_{2y} + (a_{11x} + a_{11y} + 2b_{2211})r_1 + 2b_{23}r_2], \\
c_{13} &= \frac{1}{4}[a_{23}q_{1x} - a_{23}q_{1y} + (3a_{11} + a_{33})q_{2x} - a_{3311}q_{2y} + (a_{23x} - a_{23y} - 2b_{32})q_1 + (a_{33x} - a_{33y} - 2b_{3311})q_2], \\
c_{31} &= \frac{1}{4}[3a_{32}r_{1x} - a_{32}r_{1y} + (a_{11} + 3a_{33})r_{2x} - a_{3311}r_{2y} + 2b_{32}r_1 + (a_{11x} + a_{11y} + 2b_{3311})r_2].
\end{aligned}$$

From the equation (23) we get the following system of equations:

$$q_{1t} = a_{11}q_{1xx} + b_{11}q_{1x} - c_{12x} + c_{12y} - c_{2211}q_1 - c_{32}q_2, \quad (24)$$

$$q_{2t} = a_{11}q_{2xx} + b_{11}q_{2x} - c_{13x} + c_{13y} - c_{3311}q_2 - c_{23}q_1, \quad (25)$$

$$r_{1t} = a_{22}r_{1xx} + a_{23}r_{2xx} + b_{22}r_{1x} + b_{23}r_{2x} - c_{21x} - c_{21y} + c_{2211}r_1 + c_{23}r_2, \quad (26)$$

$$r_{2t} = a_{32}r_{1xx} + a_{33}r_{2xx} + b_{32}r_{1x} + b_{33}r_{2x} - c_{31x} - c_{31y} + c_{3311}r_2 + c_{32}r_1, \quad (27)$$

$$c_{11x} + c_{21}q_1 + c_{31}q_2 + c_{12}r_1 + c_{13}r_2 + b_{12}r_{1x} + b_{13}r_{2x} - c_{11y} = 0, \quad (28)$$

$$-c_{22x} - c_{12}r_1 - c_{21}q_1 - b_{21}q_{1x} - c_{22y} = 0, \quad (29)$$

$$-c_{23x} - c_{13}r_1 - c_{21}q_2 - b_{21}q_{2x} - c_{23y} = 0, \quad (30)$$

$$-c_{32x} - c_{12}r_2 - c_{31}q_1 - b_{31}q_{1x} - c_{32y} = 0, \quad (31)$$

$$-c_{33x} - c_{13}r_2 - c_{31}q_2 - b_{31}q_{2x} - c_{33y} = 0. \quad (32)$$

Then, taking into account the above results, the equation (10) can be rewritten as

$$\begin{aligned}
iq_{1t} + \frac{i}{4}a_{2211}q_{1xx} + \frac{i}{4}a_{2211}q_{1yy} + \frac{i}{4}a_{32}q_{2xx} + \frac{i}{4}a_{32}q_{2yy} + \frac{i}{2}(-a_{22} - a_{11})q_{1xy} - \\
-\frac{i}{2}a_{32}q_{2xy} + \frac{i}{2}(a_{32x} - a_{32y} - b_{32})q_{2x} + \frac{i}{2}(-a_{32x} + a_{32y} + b_{32})q_{2y} + \\
+\frac{i}{2}(a_{22x} - a_{22y} - b_{22} - b_{11})q_{1x} + \frac{i}{2}(a_{22x} + a_{22y} + b_{22} - b_{11})q_{1y} + \\
+\frac{i}{4}[a_{22xx} - 2a_{22xy} + a_{22yy} - 2(b_{22x} - b_{22y}) + 4c_{2211}]q_1 + \\
+\frac{i}{4}[a_{32xx} - 2a_{32xy} + a_{32y} - 2(b_{32x} - b_{32y}) + 4c_{32}]q_2 = 0.
\end{aligned} \quad (33)$$

Assuming that the coefficients of the second derivative of q_1 with respect to x and y in the equation (33) are equal to unity, i. e. $\frac{i}{4}a_{2211} = 1$, we obtain $a_{2211} = a_{22} - a_{11} = -4i$. Without loss of generality, take $a_{11} = 2i$ and $a_{22} = -2i$.

Similarly, we consider the equation (25). We get that

$$\begin{aligned} iq_{2t} + \frac{i}{4}a_{3311}q_{2xx} + \frac{i}{4}a_{3311}q_{2yy} + \frac{i}{4}a_{23}q_{1xx} + \frac{i}{4}a_{23}q_{1yy} + \frac{i}{2}(-a_{33} - a_{11})q_{2xy} - \\ - \frac{i}{2}a_{23}q_{1xy} + \frac{i}{2}(a_{23x} - a_{23y} - b_{32})q_{1x} + \frac{i}{2}(-a_{23x} + a_{23y} + b_{32})q_{1y} + \\ + \frac{i}{2}(a_{33x} - a_{33y} - b_{33} - b_{11})q_{2x} + \frac{i}{2}(a_{33x} + a_{33y} + b_{33} - b_{11})q_{2y} + \\ + \frac{i}{4}[a_{33xx} - 2a_{33xy} + a_{33yy} - 2(b_{33x} - b_{33y}) + 4c_{3311}]q_2 + \\ + \frac{i}{4}[a_{23xx} - 2a_{23xy} + a_{23y} - 2(b_{23x} - b_{23y}) + 4c_{23}]q_1 = 0. \end{aligned} \quad (34)$$

We assume that in the equation $34\frac{i}{4}a_{3311} = 1$. Then we determine that $a_{33} = -2i$. Given the above, the matrix B_2 takes the form

$$B_2 = \begin{pmatrix} 2i & 0 & 0 \\ 0 & -2i & a_{23} \\ 0 & a_{32} & -2i \end{pmatrix} = 2i\Sigma + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & a_{23} \\ 0 & a_{32} & 0 \end{pmatrix}.$$

In the case when $a_{23} = a_{32} = 0$, the matrix B_2 takes the form

$$B_2 = 2i\Sigma.$$

Now we list all the equations obtained from 20 - 23 for the elements b_{ij} and c_{ij} . We have

$$\begin{aligned} b_{12} &= 2iq_1, \\ b_{13} &= 2iq_2, \\ b_{21} &= -2ir_1, \\ b_{31} &= -2ir_2, \\ b_{11x} - b_{11y} &= 0, \\ b_{22x} + b_{22y} &= 0, \\ b_{33x} + b_{33y} &= 0, \\ b_{23x} + b_{23y} &= 0, \\ b_{32x} + b_{32y} &= 0. \end{aligned}$$

The last five equations have a solution

$$b_{11} = b_{22} = b_{23} = b_{32} = b_{33} = 0.$$

Thus, for the elements of the matrix B_1 we get

$$B_1 = \begin{pmatrix} 0 & 2iq_1 & 2iq_2 \\ -2ir_1 & 0 & 0 \\ -2ir_2 & 0 & 0 \end{pmatrix} = 2iP.$$

Similarly, we define the expressions for the elements c_{ij} of the matrix B_0 in the form

$$\begin{aligned} c_{12} &= i(q_{1x} + q_{1y}), \\ c_{13} &= i(q_{2x} + q_{2y}), \\ c_{21} &= -i(r_{1x} - r_{1y}), \\ c_{31} &= -i(r_{2x} - r_{2y}). \end{aligned}$$

The remaining five elements c_{11} , c_{22} , c_{33} , c_{23} , c_{32} satisfy the following nontrivial equations:

$$c_{11x} - c_{11y} = -i(r_1 q_1 + r_2 q_2)_x - i(r_1 q_1 + r_2 q_2)_y, \quad (35)$$

$$c_{22x} + c_{22y} = i(r_1 q_1)_x - i(r_1 q_1)_y, \quad (36)$$

$$c_{33x} + c_{33y} = i(r_2 q_2)_x - i(r_2 q_2)_y, \quad (37)$$

$$c_{23x} + c_{23y} = i(r_1 q_2)_x - i(r_1 q_2)_y, \quad (38)$$

$$c_{32x} + c_{32y} = i(r_2 q_1)_x - i(r_2 q_1)_y. \quad (39)$$

Let us introduce the notations $v_1 = -ic_{2211}$, $v_2 = -ic_{3311}$, $w_1 = -ic_{32}$ and $w_2 = -ic_{23}$. Acting on v_1 with operators $D^+ = \partial_x + \partial_y$ $D^- = \partial_x - \partial_y$, we get

$$D^- D^+ (-ic_{22}) = -i(c_{22xx} - c_{22yy}),$$

$$D^+ D^- (-ic_{22}) = (r_1 q_1)_{xx} - (r_1 q_1)_{yx} - (r_1 q_1)_{xy} + i(r_1 q_1)_{yy} = (r_1 q_1)_{xx} - 2(r_1 q_1)_{xy} + i(r_1 q_1)_{yy},$$

$$D^+ D^- v_1 = (2r_1 q_1 + r_2 q_2)_{xx} + 2(r_2 q_2)_{xy} + (2r_1 q_1 + r_2 q_2)_{yy}.$$

As a result, we have

$$v_{1xx} - v_{1yy} = (2r_1 q_1 + r_2 q_2)_{xx} + 2(r_2 q_2)_{xy} + (2r_1 q_1 + r_2 q_2)_{yy},$$

which in turn gives the equation (14).

Similarly, acting by the operators $D^+ y$ D^- and v_2 , w_1 and w_2 , we obtain the following equations, respectively:

$$v_{2xx} - v_{2yy} = (r_1 q_1 + 2r_2 q_2)_{xx} + 2(r_1 q_1)_{xy} + (r_1 q_1 + 2r_2 q_2)_{yy},$$

$$w_{1xx} - w_{1yy} = (q_1 r_2)_{xx} - 2(q_1 r_2)_{xy} + (q_1 r_2)_{yy},$$

$$w_{2xx} - w_{2yy} = (q_2 r_1)_{xx} - 2(q_2 r_1)_{xy} + (q_2 r_1)_{yy}.$$

As we can see, the last three equations obtained above turn out to be the equivalent equations (15)-(17). Furthermore, using the above notation (35)-(39), we can write for equations for q_{1t} , q_{2t} , r_{1t} and r_{2t} in the form

$$\begin{aligned} i q_{1t} + q_{1xx} + q_{1yy} - v_1 q_1 - w_1 q_2 &= 0, \\ i q_{2t} + q_{2xx} + q_{2yy} - w_2 q_1 - v_2 q_2 &= 0, \\ -i r_{1t} + r_{1xx} + r_{1yy} - v_1 r_1 - w_1 r_2 &= 0, \\ -i r_{2t} + r_{2xx} + r_{2yy} - w_2 r_1 - v_2 r_2 &= 0. \end{aligned}$$

Thus, we obtained the system of equations (10)-(12), which was required to prove.

4. Conclusion

In conclusion, we note that the obtained two-component generalization of the DSI (10)-(17) equation and its Lax representation (18)-(19) are new. A detailed study of the algebraic and geometric properties of the (10)-(17) system is the subject of our further research.

The work was supported by the MES RK on the GF, contract N132, 03/12/2018.

References

- [1] Kostov N., Dandoloﬀ R., Gerdjikov V., Grahovski G. *The Manakov System as Two Moving Interacting Curves*, Proc.Int.Works. "Complex structures and vector fields", Sofia, Bulgaria. Eds.: K.Sekigawa, S.Dimiev. World Scientific (2007).
- [2] Myrzakul Akbota and Myrzakulov Ratbay. *Darboux transformations exact soliton solutions of integrable coupled spin systems related with the Manakov system*, [arXiv:1607.08151]
- [3] Myrzakul Akbota and Myrzakulov Ratbay. *Integrable geometric flows of interacting curves/surfaces, multilayer spin systems and the vector nonlinear Schrodinger equation*, [arXiv:1608.08553]
- [4] Nugmanova G., Myrzakul A. *Integrability of the two-layer spin system*, Proc.of Twentieth Int. Conf. "Geometry, Integrability and Quantization", Varna, Yoshioka (2018).
- [5] Ishimori Y. *Multi-Vortex solutions of the a two-dimensional nonlinear wave equation*, Prog. Theor. Phys. 1984. **72**, N1. pp. 33-37.
- [6] Davey A., Stewartson K. *On three-dimensional packets of surface waves*, Proc. of the Royal Society of London Series A, 1974. **338**, pp. 101–110.
- [7] Ablowitz M. A., Clarkson P. A. *Solitons, it Nonlinear Evolution Equations and Inverse Scattering* London Mathematical Society Lecture Note Series Book **149**, p. 516 ISBN 0521387302
- [8] Zhou Z., Ma W X., Zhou R. *Finite-dimensional integrable systems associated with Davey-Stewartson I equation* Nonlinearity **14** (2001) pp.701-717.